

The large deviation approach to statistical mechanics

Hugo Touchette

School of Mathematical Sciences, Queen Mary, University of London

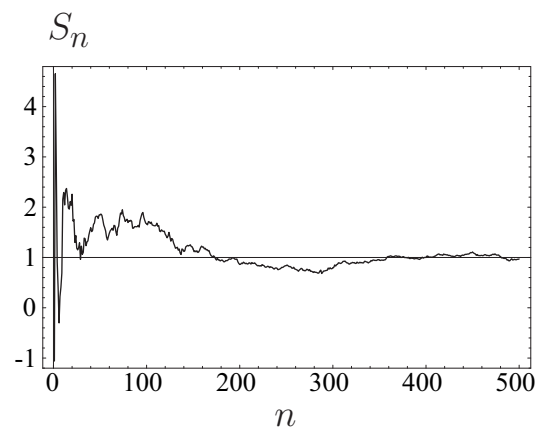
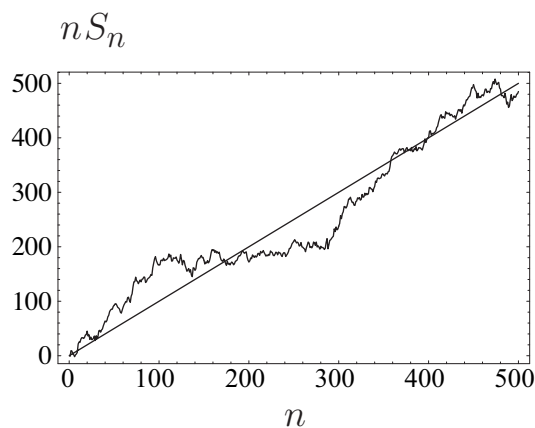
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Outline

- 1 Examples of large deviations
- 2 Basic results of large deviation theory
- 3 Mathematical applications
- 4 Physical applications

Example 1: Sum of Gaussian random variables

$$S_n = \frac{1}{n} \sum_{i=1}^n X_i, \quad X \sim N(\mu, \sigma^2)$$



Basic observations

- $S_n \rightarrow \mu$ in probability
- Fluctuations $\sim 1/\sqrt{n} \rightarrow 0$

Sum of Gaussian random variables (cont'd)

- Probability density of S_n :

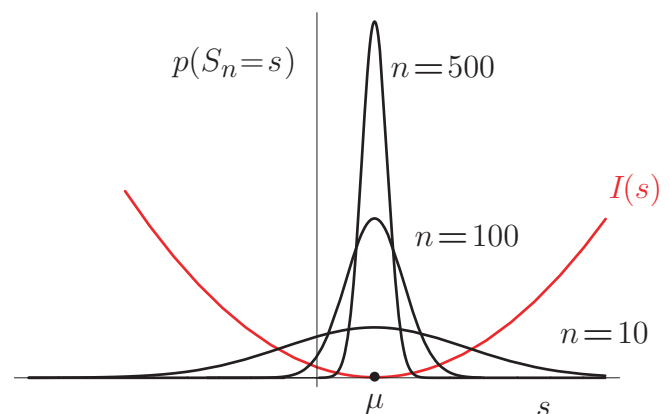
$$P(S_n = s) = \sqrt{\frac{n}{2\pi\sigma^2}} e^{-n(s-\mu)^2/(2\sigma^2)}$$

- Dominant part:

$$P(S_n = s) \approx e^{-nI(s)}$$

- Rate function:

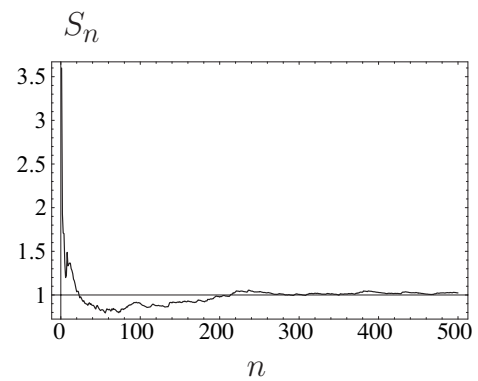
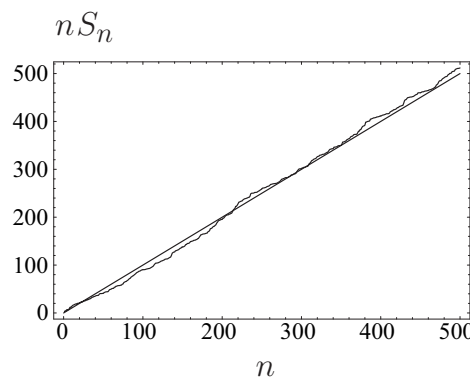
$$I(s) = \frac{(s - \mu)^2}{2\sigma^2}$$



Example 2: Sum of Poisson random variables

$$S_n = \frac{1}{n} \sum_{i=1}^n X_i$$

$$p(x) = \frac{1}{\mu} e^{-x/\mu}$$

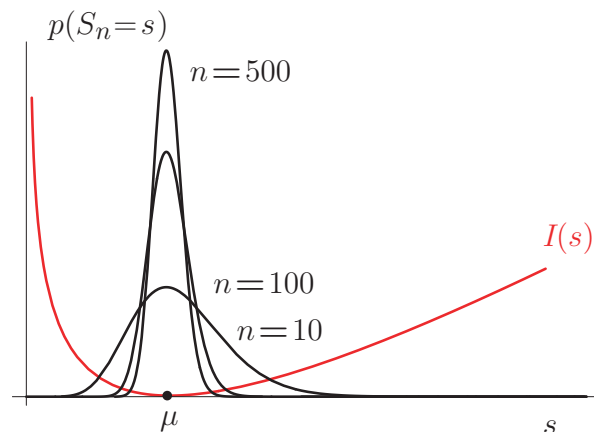


- Large deviation probability:

$$P(S_n = s) \approx e^{-nI(s)}$$

- Rate function:

$$I(s) = \frac{s}{\mu} - 1 - \ln \frac{s}{\mu}$$



Example 3: Equilibrium fluctuations

- N -particle system
- Macrostate: m

Microcanonical

Einstein (1925)

$$P_U(m) = e^{S(m,U)}$$

- Extensivity: $S \sim N$
- Large deviation probability:

$$P_U(m) \approx e^{-NI_U(m)}$$

- Equilibrium: $I_U(m^*) = 0$

Canonical

Landau (1935)

$$P_\beta(m) = e^{-F(m,\beta)}$$

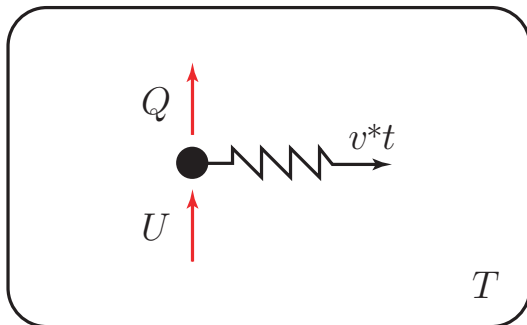
- Extensivity: $F \sim N$
- Large deviation probability:

$$P_\beta(m) \approx e^{-NI_\beta(m)}$$

- Equilibrium: $I_\beta(m^*) = 0$

Example 4: Nonequilibrium fluctuations

van Zon and Cohen (2003)



- Pulling force: $F = -k[x(t) - v^*t]$
- Langevin dynamics:

$$\dot{x}(t) = -\tau_r^{-1}[x(t) - v^*t] + \alpha^{-1}\xi(t)$$

- Work: $W_\tau = \frac{1}{\tau} \int_0^\tau Fv^* dt = \Delta U + Q_\tau$

Work fluctuations

- Large deviation probability:

$$P(W_\tau = w) \approx e^{-\tau I(w)}$$

- Fluctuation Theorem:

$$\frac{P(w)}{P(-w)} = e^{a\tau w}, \quad I(w) = I(-w) - a\tau w$$

Large deviation theory

- Random variable: A_n
- Probability density: $P(A_n = a)$

Large deviation principle

$$P(A_n = a) \approx e^{-nI(a)}$$

$$\begin{aligned} \ln P(a) &= -nI(a) + o(n) \\ \lim_{n \rightarrow \infty} -\frac{1}{n} \ln P(a) &= I(a) \end{aligned}$$

- Rate function: $I(a) \geq 0$

Goals of large deviation theory

- 1 Prove that a large deviation principle exists
- 2 Calculate the rate function

Large deviation theory (cont'd)

Large deviation principle

- Upper bound for closed sets C :

$$\limsup_{n \rightarrow \infty} \frac{1}{n} \ln P(A_n \in C) \leq - \inf_{a \in C} I(a)$$

- Lower bound for open sets O :

$$\liminf_{n \rightarrow \infty} \frac{1}{n} \ln P(A_n \in O) \geq - \inf_{a \in O} I(a)$$

- $I(a)$ is lower semicontinuous
- Bounds related to weak convergence

The Gärtner-Ellis Theorem

Log-generating function or free energy function

$$\lambda(k) = \lim_{n \rightarrow \infty} \frac{1}{n} \ln \langle e^{nkA_n} \rangle, \quad k \in \mathbb{R}$$

Theorem: Gärtner (1977), Ellis (1984)

If $\lambda(k)$ is differentiable, then

- 1 Existence of large deviation principle:

$$P(A_n = a) \approx e^{-nI(a)}$$

- 2 Expression for $I(a)$:

$$I(a) = \max_k \{ka - \lambda(k)\}$$

- $I(a)$ is the Legendre transform of $\lambda(k)$
- $I(a)$ is convex

Contraction principle

- Large deviation principle for A_n :

$$P(A_n = a) \approx e^{-nI_A(a)}$$

- Contraction: $B_n = f(A_n)$

Contraction principle

- Large deviation principle for B_n :

$$P(B_n = b) \approx e^{-nI_B(b)}$$

- Rate function:

$$I_B(b) = \inf_{a:f(a)=b} I_A(a) = \inf_{f^{-1}(b)} I_A(a)$$

Sum of IID random variables

Cramér (1938)

- Random variable:

$$S_n = \frac{1}{n} \sum_{i=1}^n X_i, \quad X_i \sim p(x), \quad IID$$

- Free energy function:

$$\lambda(k) = \lim_{n \rightarrow \infty} \frac{1}{n} \ln \langle e^{nkS_n} \rangle = \lim_{n \rightarrow \infty} \frac{1}{n} \ln \left\langle \prod_{i=1}^n e^{kX_i} \right\rangle = \ln \langle e^{kX} \rangle$$

Gaussian

$$\lambda(k) = \mu k + \frac{\sigma^2}{2} k^2, \quad k \in \mathbb{R}$$
$$I(s) = \frac{(s - \mu)^2}{2\sigma^2}, \quad s \in \mathbb{R}$$

Poisson

$$\lambda(k) = -\ln(1 - \mu k), \quad k < \frac{1}{\mu}$$
$$I(s) = \frac{s}{\mu} - 1 - \ln \frac{s}{\mu}, \quad s > 0$$

Sanov's Theorem

Sanov (1961), Boltzmann (1877)

- Sequence of IID random variables: X_1, X_2, \dots, X_n
- Empirical frequencies:

$$L_n(x) = \frac{1}{n} \sum_{i=1}^n \delta(X_i - x) = \frac{\# \text{ occurrences of } X_i = x}{n}$$

Gärtner-Ellis Theorem

- Free energy:

$$\lambda[k] = \lim_{n \rightarrow \infty} \frac{1}{n} \ln \langle e^{nk \cdot L_n} \rangle = \ln \langle e^{k(X)} \rangle$$

- Large deviation principle:

$$P(L_n = \ell) \approx e^{-nD(\ell||p)}, \quad D(\ell||p) = \int dx \ell(x) \ln \frac{\ell(x)}{p(x)}$$

Markov processes

Donsker and Varadhan (1975)

- Markov chain: $X_1 \rightarrow X_2 \rightarrow \dots \rightarrow X_n$
- Probability:

$$P(x) = \mu(x_1) \pi(x_2|x_1) \cdots \pi(x_n|x_{n-1})$$

- Additive process:

$$S_n = \frac{1}{n} \sum_{i=1}^n f(X_i)$$

Gärtner-Ellis Theorem

- Free energy function:

$$\lambda(k) = \ln \zeta(\Pi_k), \quad \pi_k(x_n|x_{n-1}) = \pi(x_n|x_{n-1}) e^{kf(x_n)}$$

- Large deviation principle:

$$P(S_n = s) \approx e^{-nI(s)}, \quad I(s) = \sup_k \{ks - \lambda(k)\}$$

Noise-perturbed dynamics

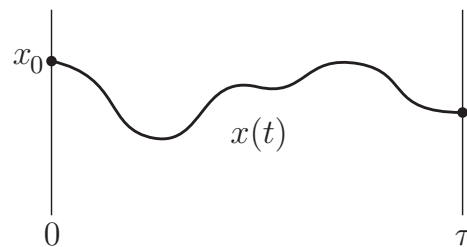
Freidlin and Wentzell (1984)

- Perturbed dynamics:

$$\dot{X}_\epsilon(t) = b(X_\epsilon) + \sqrt{\epsilon}\eta(t), \quad \epsilon \rightarrow 0$$

- Gaussian white noise:

$$\langle \eta(t) \rangle = 0, \quad \langle \eta(t)\eta(t') \rangle = \delta(t - t')$$



LDP for the paths

$$P[x] \approx e^{-J[x]/\epsilon}, \quad J[x] = \frac{1}{2} \int_0^\tau [\dot{x}(t) - b(x(t))]^2 dt$$

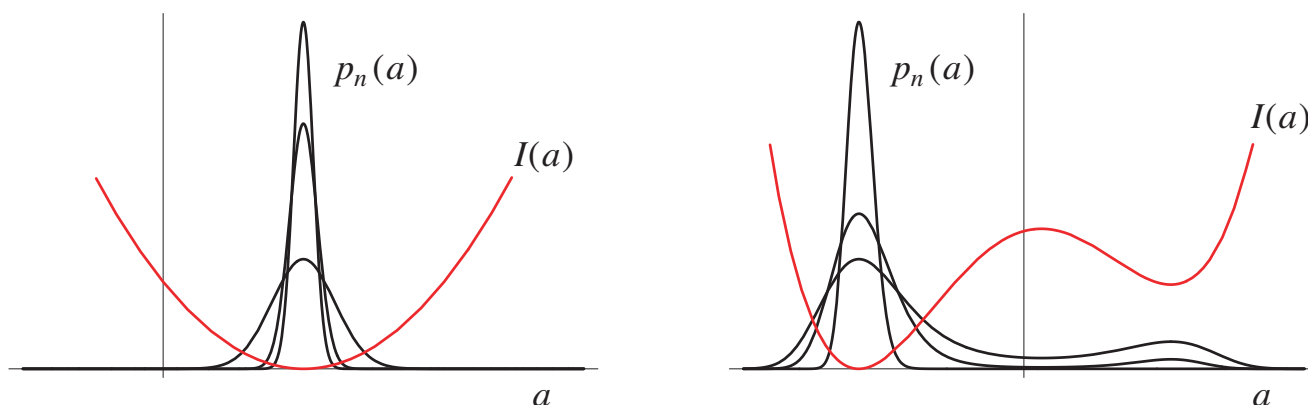
WKB approximations

$$P(x, \tau | x_0) = \int_{x(0)=x_0}^{x(\tau)=x} \mathcal{D}[x] P[x] \approx e^{-V(x, \tau | x_0)/\epsilon}$$

$$V(x, \tau | x_0) = \inf_{x(t): x(0)=x_0, x(\tau)=x} J[x]$$

Properties of I

- $I \geq 0$
- I is often convex
- Most probable value \Leftrightarrow min and zero of I
- Zero of $I \Leftrightarrow$ Law of Large Numbers
- Parabolic minimum \Leftrightarrow Central Limit Theorem



Equilibrium statistical mechanics

Lanford (1973), Ellis (1985)

- N -particle system
- Macrostate: m

Microcanonical

- Large deviation probability:

$$P_U(m) \approx e^{-NI_U(m)}$$

- Equilibrium:

$$\inf_m I_U(m) = 0$$

- Max entropy principle

Canonical

- Large deviation probability:

$$P_\beta(m) \approx e^{-NI_\beta(m)}$$

- Equilibrium:

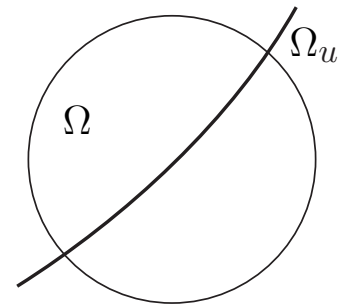
$$\inf_m I_\beta(m) = 0$$

- Min free energy principle

- Equilibrium states determined by variational principles

Microcanonical entropy

- Microstate: $\omega \in \Omega$
- Energy: $U(\omega)$
- Density of states: $\rho(u) = \int_{\Omega} d\omega \delta(U(\omega)/N - u)$
- Large deviation form: $\rho(u) \approx e^{Ns(u)}$



Gärtner-Ellis Theorem

- Free energy:

$$\varphi(\beta) = \lim_{N \rightarrow \infty} -\frac{1}{N} \ln Z(\beta), \quad Z(\beta) = \int_{\Omega} d\omega e^{-\beta U(\omega)}$$

- If $\varphi(\beta)$ is differentiable, then

$$\rho(u) \approx e^{Ns(u)} \quad s(u) = \inf_{\beta} \{\beta u - \varphi(\beta)\}$$

- Basis for Legendre transforms in statistical mechanics

Fluctuation Theorems

Farago (2002)

- Langevin dynamics:

$$\dot{v}(t) + \gamma v(t) = \underbrace{\psi(t)}_{\text{noise}}$$

- Energetics:

$$\frac{d}{dt} \left(\frac{1}{2} v^2 \right) = \underbrace{-\gamma v^2}_{\text{dissipation}} + \underbrace{\psi v}_{\text{injection}}$$

- Injected energy:

$$W_\tau = \frac{1}{\tau} \int_0^\tau \psi v dt \quad (\text{Markov additive process})$$

- Large deviation probability:

$$P(W_\tau = w) \approx e^{-\tau I(w)}$$

- ▶ Large deviations of Markov processes
- ▶ Donsker and Varadhan (1975)

Fluctuations Theorems (cont'd)

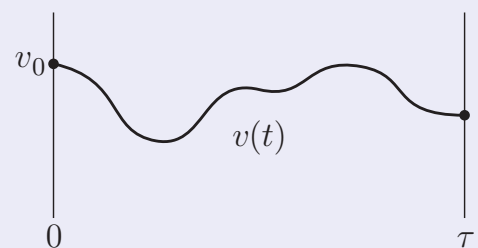
Gärtner-Ellis Theorem

- Free energy function:

$$\lambda(k) = \lim_{\tau \rightarrow \infty} \frac{1}{\tau} \ln \langle e^{\tau k W_\tau} \rangle$$

- Path integral:

$$\langle e^{\tau k W_\tau} \rangle = \int \mathcal{D}[v] P[v] e^{k \int_0^\tau v(\dot{v} + \gamma v) dt}$$



- Large deviation result:

$$P(W_\tau = w) \approx e^{-\tau I(w)}, \quad I(w) = \max_k \{kw - \lambda(k)\}$$

- Fluctuation Theorem:

$$\frac{P(w)}{P(-w)} = e^{-\tau [I(w) - I(-w)]} \neq e^{a\tau w}$$

Multifractals

- Measure: $\mu(x)$
- Coarse-grained measure:

$$p_{\varepsilon,i} = \int_{i\text{th box}} \mu(x)$$

- Structure function:

$$Z_{\varepsilon}(q) = \sum_i p_{\varepsilon,i}^q \approx \varepsilon^{\tau(q)}$$

- Multifractal spectrum:

$$n_{\varepsilon}(\alpha) \approx \varepsilon^{-f(\alpha)}$$

Gärtner-Ellis Theorem

If $\tau(q)$ is differentiable, then

$$f(\alpha) = \inf_q \{q\alpha - \tau(q)\}$$

Thermodynamic formalism

- Map: $x_{n+1} = f(x_n)$
- Invariant density: $\rho(x)$
- Time average:

$$S_n = \frac{1}{n} \sum_{i=1}^n g(x_i) \longrightarrow \langle g(x) \rangle = \int g(x) \rho(x) dx$$

- Empirical density:

$$L_n(x) = \frac{1}{n} \sum_{i=1}^n \delta(x_i - x) \longrightarrow \rho(x)$$

- Large deviation principles:

$$P(S_n = s) \approx e^{-nI(s)}, \quad P(L_n = \ell) \approx e^{-nJ(\ell)}$$

The big picture

The mathematics of statistical mechanics
is large deviation theory

- Equilibrium statistical mechanics
 - ▶ Entropy = rate function
 - ▶ Free energy = log-generating function
 - ▶ Why variational principles for equilibrium states?
 - ▶ Why the Legendre transform?
- Nonequilibrium statistical mechanics
 - ▶ Fluctuation Theorems are large deviation results
 - ▶ Variational principles for steady states
 - ▶ Entropy = rate function
- Limits of statistical mechanics = limits of large deviation theory
 - ▶ $\lambda(k)$ non-differentiable
 - ▶ Nonconvex rate function, nonconcave entropies
 - ▶ First-order phase transitions
 - ▶ Non-exponential deviations: $\lambda(k) = 0$ or ∞

Further reading



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Springer, New York, 1985.



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Further reading (cont'd)



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H. Touchette

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<http://www.maths.qmul.ac.uk/~ht>

ht@maths.qmul.ac.uk