

Anomalous dynamics of cell migration

P. Dieterich¹ R. Klages² R. Preuss³ A. Schwab⁴

¹Institute for Physiology, Dresden University of Technology

²School of Mathematical Sciences, Queen Mary University of London

³Center for Interdisciplinary Plasma Science, MPI for Plasma Physics, Garching

⁴Institute for Physiology II, University of Münster

Open Statistical Physics 2011
Open University, 2 March 2011

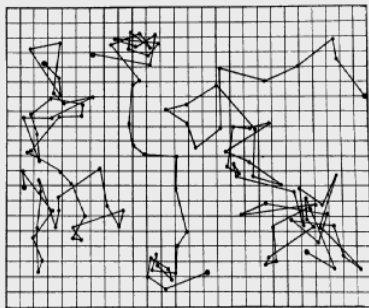


Outline

- 1 **Cell migration:** motivation and some biological details
- 2 **Experimental results:** statistics of cell migration
- 3 **Theoretical modeling:** fractional stochastic equation

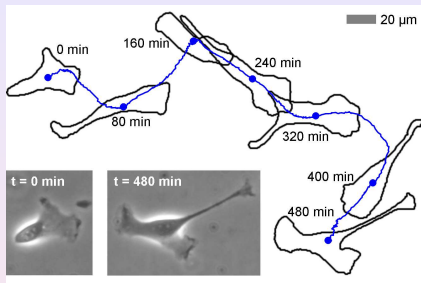
Brownian motion of migrating cells?

Brownian motion



Perrin (1913)

three colloidal particles,
positions joined by straight
lines



Dieterich et al. (2008)

single biological cell crawling on
a substrate

Brownian motion?

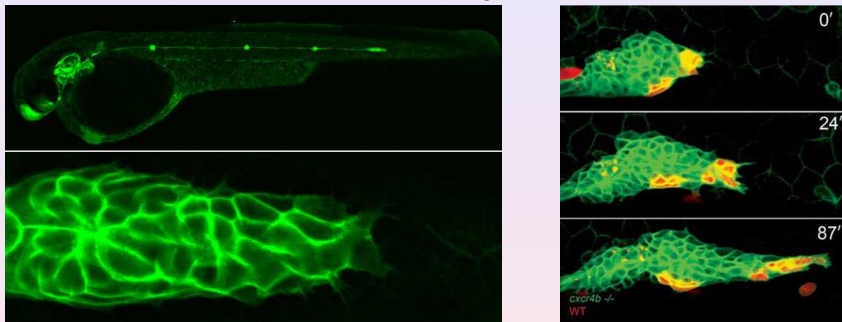
conflicting results:

yes: Dunn, Brown (1987)

no: Hartmann et al. (1994)

Why cell migration?

motion of the *primordium* in developing zebrafish:



Gilmour (2008)

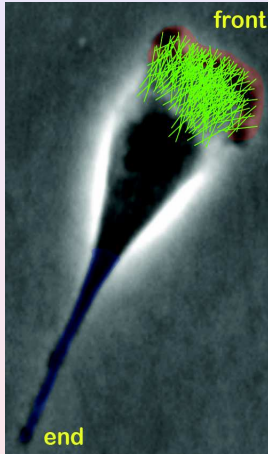
positive aspects:

- morphogenesis
- immune defense

negative aspects:

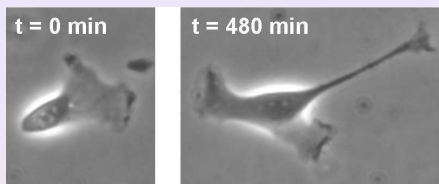
- tumor metastases
- inflammation reactions

How do cells migrate?



- **membrane protrusions and retractions** ~ force generation:
 - lamellipodia (front)
 - uropod (end)
 - actin-myosin network
- formation of a **polarized state**
front/end
- cell-substrate **adhesion**

Our cell types and some typical scales



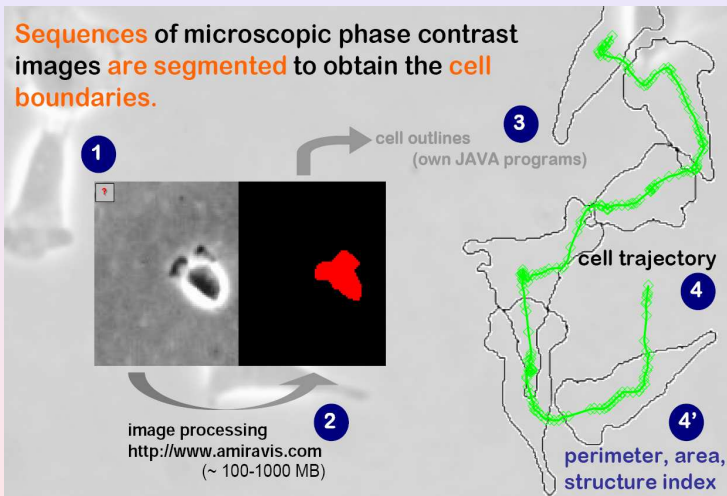
- **renal epithelial MDCK-F (Madin-Darby canine kidney) cells**; two types: wildtype (NHE^+) and NHE -deficient (NHE^-)
- observed up to **1000 minutes**: here *no* limit $t \rightarrow \infty$!
- cell diameter **$20\text{-}50\mu\text{m}$** ; mean velocity $\sim 1\mu\text{m}/\text{min}$; lamellipodial dynamics \sim **seconds**

movies: NHE+: t=210min, dt=3min

NHE-: t=171min, dt=1min

Measuring cell migration

Sequences of microscopic phase contrast images **are segmented** to obtain the **cell boundaries**.



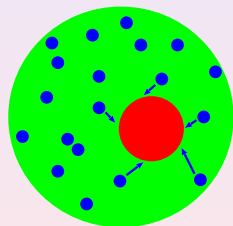
Theoretical modeling of Brownian motion

'Newton's law of stochastic physics':

$$\dot{\mathbf{v}} = -\kappa\mathbf{v} + \sqrt{\zeta}\xi(t) \quad \text{Langevin equation (1908)}$$

for a **tracer particle of velocity \mathbf{v}** immersed in a fluid

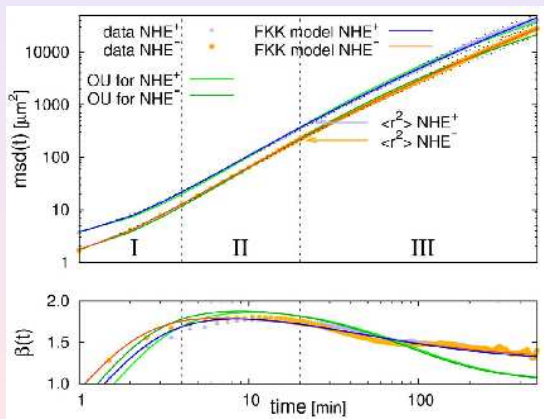
force decomposed into **viscous damping** and **random kicks of surrounding particles**



Application to cell migration?

Mean square displacement

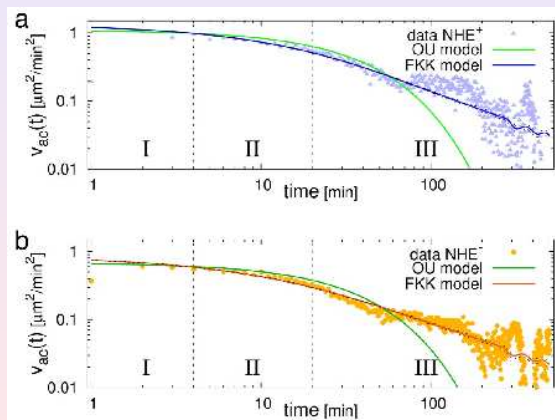
- $msd(t) := \langle [\mathbf{x}(t) - \mathbf{x}(0)]^2 \rangle \sim t^\beta$ with $\beta \rightarrow 2$ ($t \rightarrow 0$) and $\beta \rightarrow 1$ ($t \rightarrow \infty$) for Brownian motion; $\beta(t) = d \ln msd(t) / d \ln t$



anomalous diffusion if $\beta \neq 1$ ($t \rightarrow \infty$); here: **superdiffusion**

Velocity autocorrelation function

- $v_{ac}(t) := \langle \mathbf{v}(t) \cdot \mathbf{v}(0) \rangle \sim \exp(-\kappa t)$ for Brownian motion
- fits with same parameter values as $msd(t)$



crossover from **stretched exponential to power law**

Position distribution function

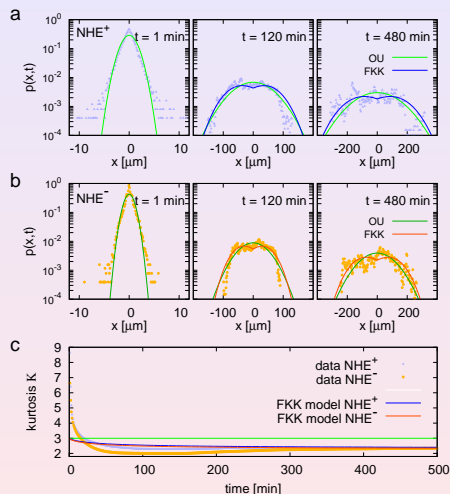
- $P(x, t) \rightarrow$ Gaussian ($t \rightarrow \infty$) and kurtosis

$$\kappa(t) := \frac{\langle x^4(t) \rangle}{\langle x^2(t) \rangle^2} \rightarrow 3 \quad (t \rightarrow \infty)$$

for Brownian motion (green lines, in 1d)

- other solid lines: fits from our model; parameter values as before

note: model needs to be amended to explain short-time distributions



crossover from peaked to broad **non-Gaussian distributions**

The model

Fractional Klein-Kramers equation (Barkai, Silbey, 2000):

$$\frac{\partial P}{\partial t} = -\frac{\partial}{\partial x} [vP] + \frac{\partial^{1-\alpha}}{\partial t^{1-\alpha}} \kappa \left[\frac{\partial}{\partial v} v + v_{th}^2 \frac{\partial^2}{\partial v^2} \right] P$$

with probability distribution $P = P(x, v, t)$, damping term κ , thermal velocity v_{th} and **Riemann-Liouville fractional derivative of order $1 - \alpha$** defined by

$$\frac{\partial^\gamma P}{\partial t^\gamma} := \begin{cases} \frac{\partial^m P}{\partial t^m} & , \quad \gamma = m \\ \frac{\partial^m}{\partial t^m} \left[\frac{1}{\Gamma(m-\gamma)} \int_0^t dt' \frac{P(t')}{(t-t')^{\gamma+1-m}} \right] & , \quad m-1 < \gamma < m \end{cases}$$

with $m \in \mathbb{N}$; for $\alpha = 1$ ordinary Klein-Kramers equation recovered

4 fit parameters v_{th}, α, κ (plus another one for ‘biological noise’ on short time scales)

Solutions for this model

analytical solutions (Barkai, Silbey, 2000):

- **mean square displacement:**

$$msd(t) = 2v_{th}^2 t^2 E_{\alpha,3}(-\kappa t^\alpha) \rightarrow 2 \frac{D_\alpha t^{2-\alpha}}{\Gamma(3-\alpha)} \quad (t \rightarrow \infty)$$

with $D_\alpha = v_{th}^2 / \kappa$ and *generalized Mittag-Leffler function*

$$E_{\alpha,\beta}(z) = \sum_{k=0}^{\infty} \frac{z^k}{\Gamma(\alpha k + \beta)}, \quad \alpha, \beta > 0, \quad z \in \mathbb{C};$$

note that $E_{1,1}(z) = \exp(z)$: $E_{\alpha,\beta}(z)$ is a generalized exponential function

- **velocity autocorrelation function:**

$$v_{ac}(t) = v_{th}^2 E_{\alpha,1}(-\kappa t^\alpha) \rightarrow \frac{1}{\kappa \Gamma(1-\alpha) t^\alpha} \quad (t \rightarrow \infty)$$

- for $\kappa \rightarrow \infty$ fractional Klein-Kramers reduces to a *fractional diffusion equation* yielding $P(x, t)$ in terms of a Fox function (Schneider, Wyss, 1989)

Possible physical interpretation

Physical meaning of the fractional derivative?

fractional Klein-Kramers equation is *approximately* related to the generalized Langevin equation

$$\dot{v} + \int_0^t dt' \kappa(t-t')v(t') = \sqrt{\zeta} \xi(t)$$

e.g., Mori, Kubo (1965/66)

with **time-dependent friction coefficient** $\kappa(t) \sim t^{-\alpha}$

cell anomalies might originate from **glassy behavior** of the cytoskeleton gel, where power law exponents are conjectured to be universal (Fabry et al., 2003; Kroy et al., 2008)

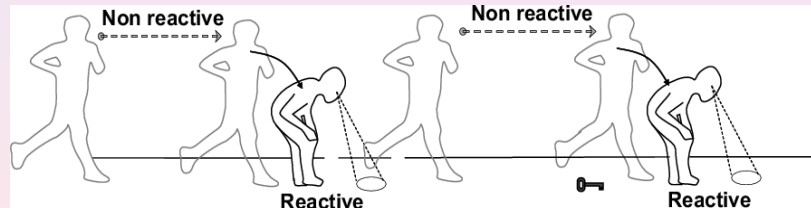
note: anomalous dynamics observed for *6 different cell types*

Possible biological interpretation

Biological meaning of the anomalous cell migration?

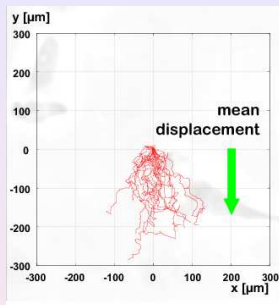
experimental data and theoretical modeling suggest *slower diffusion for small times* while *long-time motion is faster*

compare with **intermittent optimal search strategies** of foraging animals (Bénichou et al., 2006)



note: controversy about **modeling the migration of foraging animals** (albatros, bumblebees, fruitflies,...)

Outlook: cell migration under chemical gradients



new experiments on **murine neutrophils** under **chemotaxis**:

- **linear drift** in the direction of the gradient, $\langle y(t) \rangle \sim t$
 - $msd(t) - \langle y(t) \rangle^2 \sim t^\beta$ with same exponent β as in equilibrium
- \Rightarrow **fluctuation dissipation relation**

modeled by the **fractional Klein-Kramers equation** with external force $F(x)$

$$\frac{\partial P}{\partial t} = -\frac{\partial}{\partial x} [vP] + \frac{\partial^{1-\alpha}}{\partial t^{1-\alpha}} \kappa \left[\frac{\partial}{\partial v} v - \frac{F}{\kappa m} \frac{\partial}{\partial v} + v_{th}^2 \frac{\partial^2}{\partial v^2} \right] P$$

Metzler, Sokolov (2002)

Thanks and literature

- **thanks** to A.V.Checkkin and E.Lutz for helpful discussions.
- **reference to this talk:**

P.Dieterich, R.K., R.Preuss, A.Schwab, *Anomalous Dynamics of Cell Migration*, PNAS **105**, 459 (2008)

- **as a general reference:**

R.K., G.Radons, I.M.Sokolov (Eds.)
Anomalous transport
 (Wiley-VCH, 2008)

