Outline	Normal FRs	Anomalous TFRs	Experiments	Summary
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Anomalous Fluctuation Relations

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Outline				

• 'Normal' fluctuation relations:

motivation and warm-up

• Anomalous fluctuation relations:

check transient fluctuation relations for correlated Gaussian stochastic dynamics

Relations to experiments:

glassy dynamics and cell migration

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Consider a (classical) particle system evolving from some initial state into a nonequilibrium steady state.

Measure the probability distribution $\rho(\xi_t)$ of entropy production

 ξ_t during time *t*:

$$\ln \frac{\rho(\xi_t)}{\rho(-\xi_t)} = \xi_t$$

Transient Fluctuation Relation (TFR)

Evans, Cohen, Morriss (1993); Gallavotti, Cohen (1995)

why important? of very general validity and

- generalizes the Second Law to small noneq. systems
- connection with fluctuation dissipation relations
- can be checked in experiments (Wang et al., 2002)



meaning of TFR in terms of the Second Law:



 $\rho(\xi_t) = \rho(-\xi_t) \exp(\xi_t) \ge \rho(-\xi_t) \ (\xi_t \ge \mathbf{0}) \ \Rightarrow <\xi_t > \ge \mathbf{0}$

sample specifically the tails of the pdf (large deviation result)

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warmup: check TFR for the overdamped Langevin equation

 $\dot{\mathbf{x}} = \mathbf{F} + \zeta(t)$ (set all irrelevant constants to 1)

with constant field *F* and Gaussian white noise $\zeta(t)$.

entropy production ξ_t is equal to (mechanical) work $W_t = Fx(t)$ with $\rho(W_t) = F^{-1}\varrho(x, t)$; remains to solve corresponding Fokker-Planck equation for initial condition x(0) = 0:

the position pdf is Gaussian,

$$\varrho(\mathbf{x},t) = \frac{1}{\sqrt{2\pi\sigma_{\mathbf{x}}^2}} \exp\left(-\frac{(x-\langle \mathbf{x} \rangle)^2}{2\sigma_{\mathbf{x}}^2}\right)$$

straightforward:

(work) TFR holds if
$$< x > = \sigma_x^2/2$$

and \exists fluctuation-dissipation relation 1 (FDR1) \Rightarrow TFR

see, e.g., van Zon, Cohen, PRE (2003)



goal: check TFR for Gaussian stochastic processes defined by the (overdamped) generalized Langevin equation

$$\int_{0}^{t} dt' \dot{x}(t') \mathcal{K}(t-t') = \mathcal{F} + \zeta(t)$$

e.g., Kubo (1965)

with Gaussian noise $\zeta(t)$ and memory kernel K(t)

such dynamics can generate anomalous diffusion: $\sigma_x^2 \sim t^{\alpha}$ with $\alpha \neq 1 \ (t \to \infty)$

examples of applications: polymer dynamics (Panja, 2010); biological cell migration (Dieterich et al., 2008)



consider two generic cases:

1. internal Gaussian noise defined by the FDR2,

 $<\zeta(t)\zeta(t')>\sim K(t-t')$,

with non-Markovian (correlated) noise; e.g., $K(t) \sim t^{-\beta}$

solving the corresponding generalized Langevin equation in Laplace space yields $FDR2 \Rightarrow FDR1'$

and since $\rho(W_t) \sim \varrho(x, t)$ is Gaussian

 $`\mathsf{FDR1'} \Rightarrow \mathsf{TFR}$

for correlated internal Gaussian noise \exists TFR



2. external Gaussian noise for which there is no FDR2, modeled by the (overdamped) generalized Langevin equation

 $\dot{\boldsymbol{x}} = \boldsymbol{F} + \zeta(\boldsymbol{t})$

consider two types of Gaussian noise correlated by $g(\tau) = \langle \zeta(t)\zeta(t') \rangle_{\tau=t-t'} \sim (\Delta/\tau)^{\beta}$ for $\tau > \Delta$, $\beta > 0$:



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TFRs f	or correlated	external Gau	ssian noise I	
persi	stent noise:			
result	ts for σ_x^2 and the f	fluctuation ratio R	$R(W_t) = \ln \frac{\rho(W_t)}{\rho(-W_t)}$	
• 0 <	β<1:			
supe	rdiffusion $\sigma_{\rm X}^2 \sim t^2$	$^{-\beta}$ with anomalou	is TFR $R \sim rac{VV_t}{t^{1-eta}}$	
• $\beta =$: 1:			
weak	superdiffusion σ	$\frac{2}{x} \sim t \ln\left(\frac{t}{\Delta}\right)$ with	weakly anomalous	TFR
R \sim	$W_t/\ln\left(\frac{t}{\Delta}\right)$			
• 1 <	$\beta < \infty$:			
norm	al diffusion $\sigma_x^2 \sim$	2 <i>Dt</i> with $D = \int_0^\infty$	$d\tau g(\tau)$ and anoma	lous
(gene	eralized) TFR R ~	$\frac{v v_t}{D}$		



TFRs for correlated external Gaussian noise II

antipersistent noise:

 $\int_0^\infty d\tau g(\tau) > 0 \text{ yields normal diffusion with a generalized TFR}$ for $t \gg \Delta$; for 'pure' antipersistent case with $\int_0^\infty d\tau g(\tau) = 0$:

- The regime $0 < \beta < 1$ does not exist (spectral density <0)
- 1 < β < 2: subdiffusion $\sigma_x^2 \sim t^{2-\beta}$ with anomalous TFR $R \sim W_t t^{\beta-1}$
- $\beta = 2$: weak subdiffusion $\sigma_x^2 \sim \ln(t/\Delta)$ with anomalous TFR $R \sim W_t t / \ln(t/\Delta)$
- 2 < β < ∞ : localization $\sigma_x^2 = const$. with anomalous TFR $R \sim W_t t$



relation between TFR and FDR I,II for correlated Gaussian stochastic dynamics: ('normal FR'= conventional TFR)





Relations to experiments: A glassy lattice gas

$$R(W_t) = \ln \frac{\rho(W_t)}{\rho(-W_t)} = \mathbf{f}_{\beta}(\mathbf{t}) W_t$$

means by plotting *R* for different *t* the slope should change. example 1: 2^{5}

computer simulations for glassy lattice gas with external field *E*

Sellitto, PRE (2009)



similar simulation results for three other models exhibiting glassy dynamics: Crisanti et al., PRL (2013)

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Biological cell migration					

example 2:

single biological cell crawling on a substrate; trajectory recorded with a video camera (Dieterich et al., 2008)





experiments on murine neutrophils under chemotaxis:



Dieterich et al. (2013)

- linear drift in the direction of the gradient, $\langle x(t) \rangle \sim t$
- $\sigma_x^2 \sim t^\beta$ with $\beta > 1$ (long *t*): $\not\exists$ FDR1
- some relation to the generalized Langevin equation with external noise and $0 < \beta < 1$ discussed before

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Summary				

- TFR tested for two generic cases of correlated Gaussian stochastic dynamics:
 - internal noise: FDR2 implies the validity of the 'normal' work TFR
 external noise: FDR2 is broken; sub-classes of persistent and anti-persistent noise yield both anomalous TFRs
- anomalous TFRs appear to be important for glassy dynamics: cf. computer simulations on various glassy models and experiments on ('gelly') cell migration

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Reference	es			

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